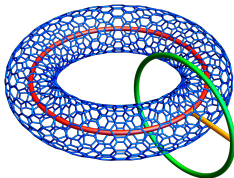


Generalized string nets at finite temperature

Julien Vidal

Laboratoire de Physique Théorique de la Matière Condensée
CNRS, Sorbonne Université, Paris



in collaboration with:

- Jean-Noël Fuchs and Anna Ritz-Zwilling (LPTMC, Paris)
- Steven H. Simon (Rudolf Peierls Centre for Theoretical Physics, Oxford)

J. Vidal, *Phys. Rev. B* **105**, L041110 (2022) / [arXiv:2108.13425](https://arxiv.org/abs/2108.13425)

A. Ritz-Zwilling, J.-N. Fuchs, S. H. Simon, and J. Vidal, [arXiv:2309.00343](https://arxiv.org/abs/2309.00343)

Topological quantum order in condensed matter in three dates

- 1989 : High- T_c superconductors, FQHE (X.-G. Wen, F. Wilczek, A. Zee)
- 1997 : Fault-tolerant quantum computation (A. Kitaev, J. Preskill)
- 2005 : String-net condensation (M. Levin, X.-G. Wen)

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Concepts coming from 70's and 80's

- Lattice gauge theories (Wegner, Wilson, Kogut)
- Conformal field theories (Pasquier, Verlinde, Moore, Seiberg)
- Topological quantum field theories (Witten)
- Knot theory (Kauffman, Jones)

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In this talk:

Two-dimensional gapped quantum systems with:

- Topology-dependent ground-state degeneracy
- Anyonic excitations

Outline

- 1 The Levin-Wen model in a nutshell
- 2 Spectral degeneracies
- 3 Partition function

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The Levin-Wen model in a nutshell

The Levin-Wen (string-net) model

- Microscopic lattice model

M. Levin and X.-G. Wen, *Phys. Rev. B* **71**, 045110 (2005)

- Trivalent graph (honeycomb lattice, two-leg ladder,...)

The Levin-Wen model in a nutshell

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- Microscopic lattice model M. Levin and X.-G. Wen, Phys. Rev. B **71**, 045110 (2005)
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Input: Unitary Fusion Category \mathcal{C}

- Set of objects (strings, labels, charges, superselection sectors, particles...)
- Fusion rules: $a \times b = \sum_c N_{ab}^c c \rightarrow$ Quantum dimensions
- Associativity of the fusion rules $\rightarrow F$ -symbols

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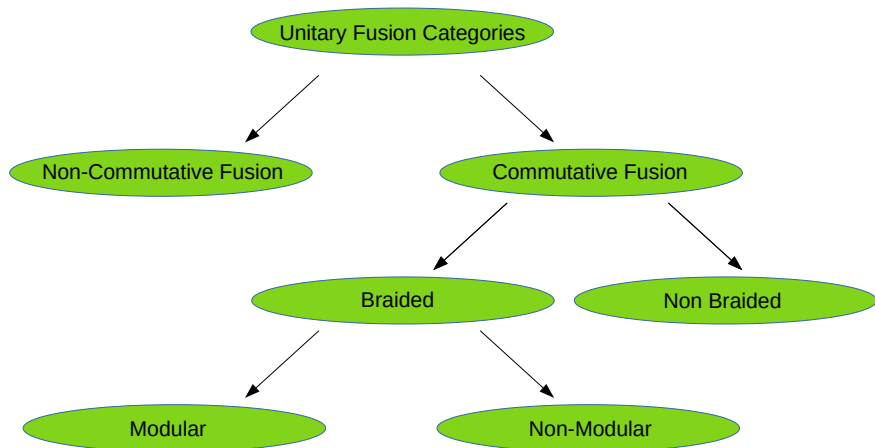
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Output: Unitary Modular Tensor Category $\mathcal{Z}(\mathcal{C})$ (Drinfeld center of \mathcal{C})

- Anyon theory (other objects, other fusion rules, other...)
- Achiral topological phase (chiral central charge $c \bmod 8 = 0$)
- Different \mathcal{C} may have the same $\mathcal{Z}(\mathcal{C})$ (Morita equivalence)

The Levin-Wen model in a nutshell

A possible tree of UFC



The Levin-Wen model in a nutshell

Local constraints and Hilbert space

- Degrees of freedom are the objects of \mathcal{C}
- Degrees of freedom are defined on the links of the trivalent graph
- Hilbert space \mathcal{H} : set of configurations respecting fusion rules at each vertex

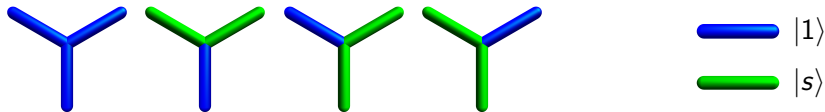
The Levin-Wen model in a nutshell

Local constraints and Hilbert space

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- Hilbert space \mathcal{H} : set of configurations respecting fusion rules at each vertex

Ex: \mathbb{Z}_2 category

- Two labels: $\{1, s\}$
- Fusion rules: $1 \times 1 = 1$, $1 \times s = s$, $s \times s = 1$



Hilbert space dimension for any trivalent graph with N_v trivalent vertices

- $\text{Dim } \mathcal{H} = 2^{\frac{N_v}{2} + 1}$

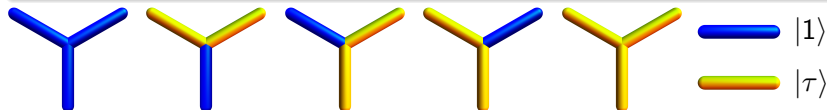
The Levin-Wen model in a nutshell

Local constraints and Hilbert space

- Degrees of freedom are the objects of \mathcal{C}
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- Hilbert space \mathcal{H} : set of configurations respecting fusion rules at each vertex

Ex: Fibonacci category

- Two objects: $\{1, \tau\}$
- Fusion rules: $1 \times 1 = 1$, $1 \times \tau = \tau$, $\tau \times \tau = 1 + \tau$



Hilbert space dimension for any trivalent graph with N_v vertices

- $\dim \mathcal{H} = (1 + \varphi^2)^{\frac{N_v}{2}} + (1 + \varphi^{-2})^{\frac{N_v}{2}}$, $\varphi = \frac{1 + \sqrt{5}}{2}$ (golden ratio)

The Levin-Wen model in a nutshell

The Hamiltonian

$$H = - \sum_p P_p$$

- H is a sum of local commuting projectors: $[P_p, P_{p'}] = 0$
- P_p : projector onto the vacuum of $\mathcal{Z}(\mathcal{C})$ in the plaquette p

$$P_p \begin{array}{c} a \quad \alpha \quad b \\ \zeta \quad \beta \quad c \\ e \quad \delta \quad d \end{array} = \sum_s \frac{d_s}{D^2} \sum' F_{s\alpha'\zeta'}^{a\zeta\alpha} F_{s\beta'\alpha'}^{b\alpha\beta} F_{s\gamma'\beta'}^{c\beta\gamma} F_{s\delta'\gamma'}^{d\gamma\delta} F_{s\epsilon'\delta'}^{e\delta\epsilon} F_{s\zeta'\epsilon'}^{f\epsilon\zeta} \begin{array}{c} a \quad \alpha' \quad b \\ \zeta' \quad \beta' \quad c \\ e \quad \delta' \quad d \end{array}$$

- d_s : quantum dimension of the string $s \in \mathcal{C}$
- $D = (\sum_s d_s^2)^{1/2}$: total quantum dimension of \mathcal{C}

Outline

- 1 The Levin-Wen model in a nutshell
- 2 Spectral degeneracies
- 3 Partition function

Spectral degeneracies

On a genus- g orientable compact surface (sphere, torus,...)

- Ground-state degeneracy: $\mathcal{D}_0 = \sum_{i \in \mathcal{Z}(\mathcal{C})} \left(\frac{d_i}{D}\right)^x = \text{Turaev-Viro invariant}$
- D : total quantum dimension of $\mathcal{Z}(\mathcal{C})$
- Euler-Poincaré characteristic: $\chi = 2 - 2g$
- $g = 0$: $\mathcal{D}_0 = 1$
- $g = 1$: $\mathcal{D}_0 = \text{Number of objects in } \mathcal{Z}(\mathcal{C})$
- $g \geq 2$: \mathcal{D}_0 depends (non trivially) on $\mathcal{Z}(\mathcal{C})$

Z. Kádár, A. Marzuoli, and M. Rasetti, Adv. Math. Phys. **2010**, 671039 (2010)

F. J. Burnell and S. H. Simon, Ann. Phys. **325**, 2550 (2010)

V. G. Turaev and O. Y. Viro, Topology **31**, 865 (1992)

G. Moore and N. Seiberg, Comm. Math. Phys. **123**, 177 (1989)

E. Verlinde, Nucl. Phys. B **300**, 360 (1988)

Spectral degeneracies

Ex: Fibonacci category (commutative, braided, modular)

- Objects of \mathcal{C} : $\{1, \tau\}$
- $d_1 = 1, d_\tau = \frac{1+\sqrt{5}}{2}$
- Objects of $\mathcal{Z}(\mathcal{C})$: $\{(1, 1), (1, \tau), (\tau, 1), (\tau, \tau)\}$
- $d_{(i,j)} = d_i d_j$

M. D. Schulz, S. Dusuel, K. P. Schmidt, and J. Vidal, Phys. Rev. Lett. **110**, 147203 (2013)

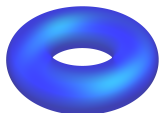
Y. Hu, S. D. Stirling, and Y.-S. Wu, Phys. Rev. B **85**, 075107 (2012)

$g = 0$



$\mathcal{D}_0 = 1$

$g = 1$



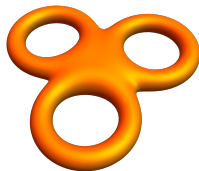
$\mathcal{D}_0 = 4$

$g = 2$



$\mathcal{D}_0 = 25$

$g = 3$



$\mathcal{D}_0 = 225$

Spectral degeneracies

Ex: $\text{Rep}(S_3)$ category (commutative, braided, non-modular)

- Objects of \mathcal{C} : $\{1, 2, 3\}$
- $d_1 = 1, d_2 = 1, d_3 = 2$
- Objects of $\mathcal{Z}(\mathcal{C})$: $\{A, B, C, D, E, F, G, H\}$
- $d_A = 1, d_B = 1, d_C = 2, d_D = 3, d_E = 3, d_F = 2, d_G = 2, d_H = 2$

A. Kitaev, *Ann. Phys.* **303**, 2 (2003)

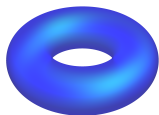
S. Beigi, P. W. Shor, D. Whalen, *Commun. Math. Phys.*, **306**, 663 (2011)

$g = 0$



$\mathcal{D}_0 = 1$

$g = 1$



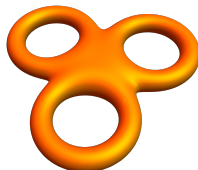
$\mathcal{D}_0 = 8$

$g = 2$



$\mathcal{D}_0 = 116$

$g = 3$



$\mathcal{D}_0 = 2948$

Spectral degeneracies

Ex: Tambara-Yamagami category TY_3 (commutative, non-braided)

- Objects of \mathcal{C} : $\{1, 2, 3, \sigma\}$
- $d_1 = d_2 = d_3 = 1, d_\sigma = \sqrt{3}$
- Objects of $\mathcal{Z}(\mathcal{C})$: $\{\alpha_{i=1, \dots, 15}\}$
- $d_{i=1, \dots, 6} = 1, d_{i=7, \dots, 9} = 2, d_{i=10, \dots, 15} = \sqrt{3}$

D. Tambara and S. Yamagami, *J. Algebra* **209**, 692 (1998)

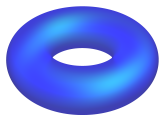
S. Gelaki, D. Naidu, and D. Nikshych, *Algebra Number Theory* **3**, 959 (2009)

$g = 0$



$\mathcal{D}_0 = 1$

$g = 1$



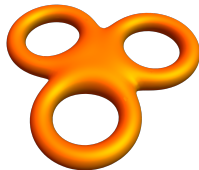
$\mathcal{D}_0 = 15$

$g = 2$



$\mathcal{D}_0 = 315$

$g = 3$



$\mathcal{D}_0 = 8883$

Spectral degeneracies

Ex: Haagerup subfactor category \mathcal{H}_3 (non-commutative)

- Objects of \mathcal{C} : $\{1, \alpha, \alpha^*, \rho, \rho\alpha, \rho\alpha^*\}$
- $d_1 = d_\alpha = d_{\alpha^*} = 1, d_\rho = d_{\rho\alpha} = d_{\rho\alpha^*} = \frac{3+\sqrt{13}}{2}$
- Objects of $\mathcal{Z}(\mathcal{C})$: $\{0, \mu^1, \mu^2, \mu^3, \mu^4, \mu^5, \mu^6, \pi^1, \pi^2, \sigma^1, \sigma^2, \sigma^3\}$
- $d_0 = 1, d_\mu = 3d_\rho, d_\pi = 3d_\rho + 1, d_\sigma = 3d_\rho + 2$

M. Asaeda and U. Haagerup, *Comm. Math. Phys.* **202**, 1 (1999)

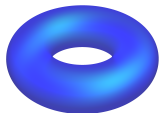
S.-M. Hong, E. Rowell, and Z. Wang, *Commun. Contemp. Math.* **10**, 1049 (2008)

$g = 0$



$\mathcal{D}_0 = 1$

$g = 1$



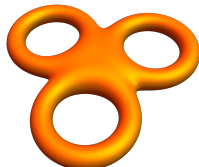
$\mathcal{D}_0 = 12$

$g = 2$



$\mathcal{D}_0 = 1401$

$g = 3$



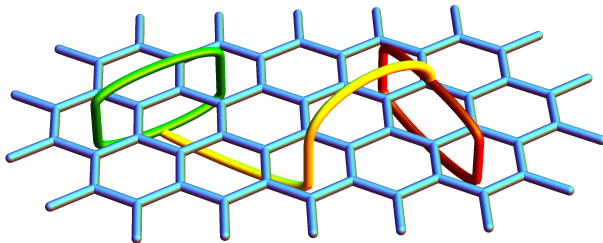
$\mathcal{D}_0 = 1\,603\,329$

Spectral degeneracies

Excitations of the Levin-Wen model (2005)

- Excitations are the nontrivial objects of $\mathcal{Z}(\mathcal{C})$
- General construction of the Drinfeld center via the Ocneanu tube algebra
- Some objects only violate the plaquette constraints (fluxons)
- Some objects violate the vertex/plaquette constraints (not considered here !)

T. Lan and X.-G. Wen, Phys. Rev. B **90**, 115119 (2014)



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Old results for commutative \mathcal{C}

- Number of states with q fluxons on a genus- g surface with N_p plaquettes:

$$\mathcal{D}_q = \binom{N_p}{q} \sum_{i \in \mathcal{Z}(\mathcal{C})} \left(\frac{d_i}{D} \right)^{\chi - q} \left(1 - \frac{d_i}{D} \right)^q$$

G. Moore and N. Seiberg, *Comm. Math. Phys.* **123**, 177 (1989)

J. Vidal, *Phys. Rev. B* **105**, L041110 (2022)

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- n_i : internal multiplicity of the particle i given by the tube algebra

C.-H. Lin, M. Levin, and F. J. Burnell, *Phys. Rev. B* **103**, 195155 (2021)

A. Ritz-Zwilling, J.-N. Fuchs, S. H. Simon, and J. Vidal, [arXiv:2309.00343](https://arxiv.org/abs/2309.00343)

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A. Ritz-Zwilling, J.-N. Fuchs, S. H. Simon, and J. Vidal, arXiv:2309.00343

Example of Morita equivalence: $\text{Rep}(S_3)$ and $\text{Vec}(S_3)$

- Same Drinfeld center $\mathcal{Z}(S_3)$: $\{A, B, C, D, E, F, G, H\}$
- $\text{Rep}(S_3)$: $\mathbf{n} = (1, 0, 0, 1, 0, 1, 0, 0)$
- $\text{Vec}(S_3)$: $\mathbf{n} = (1, 1, 2, 0, 0, 0, 0, 0)$
- Different Hilbert space, different degeneracies,...

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Partition function

Energy spectrum

- Levin-Wen Hamiltonian: $H = - \sum_p P_p$
- Ground-state energy: $E_0 = -N_p$
- q -fluxon state energy: $E_q = -N_p + q$

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Exact finite-size and finite-temperature partition function

- Partition function: $Z = \text{Tr}(e^{-\beta H}) = \sum_{q=0}^{N_p} \mathcal{D}_q e^{-\beta E_q} \quad \left(\beta = \frac{1}{k_B T}\right)$

$$Z = \sum_{i \in \mathcal{Z}(\mathcal{C})} \left(\frac{d_i}{D}\right)^{\chi - N_p} \left[n_i - \left(\frac{d_i}{D}\right) (1 - e^\beta) \right]^{N_p}$$

- Z depends on the fusion rules: $d_a \times d_b = \sum_c N_c^{ab} d_c$
- Z depends on the surface topology: $\chi = 2 - 2g$
- Z depends on the number of plaquettes: N_p

Partition function

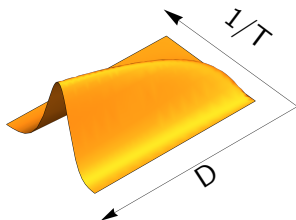
Specific heat and absence of thermal phase transition

- Specific heat per plaquette : $c = \frac{\beta^2}{N_p} \frac{\partial^2 \ln Z}{\partial \beta^2}$
- Thermodynamical limit ($N_p \rightarrow \infty$): $c = \frac{e^\beta \beta^2 (D-1)}{(D-1 + e^\beta)^2}$
- Only depends on the total quantum dimension D

No finite-temperature phase transition in the Levin-Wen model !

J. Vidal, Phys. Rev. B **105**, L041110 (2022)

A. Ritz-Zwilling, J.-N. Fuchs, S. H. Simon, and J. Vidal, in preparation



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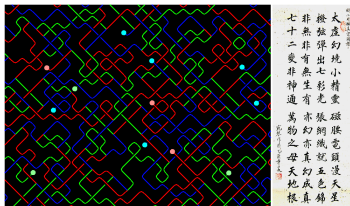
Absence of topological order at $T > 0$ for any two-dimensional Hamiltonian which is a sum of local commuting projectors !

M. Hastings, Phys. Rev. Lett. **107**, 210501 (2011)

Outlook

Summary and extensions

- Exact partition function of the Levin-Wen model
- More results available for:
 - any UFC with fusion multiplicities ($N_{ab}^c \geq 1$)
 - any topology with punctures (boundaries)
- Absence of finite-temperature phase transition
- Observables (e.g., Wegner-Wilson loops)
- Importance of charge excitations (violation of Gauss's law)



Courtesy of X.-G.Wen

What's next ?

Forthcoming conference in Les Houches (french Alps): 1-12 April 2024

Topological order: Anyons and Fractons

Website: <https://topoanyons.sciencesconf.org/>